

ON CUBATURE RULES ASSOCIATED TO WEYL GROUP ORBIT FUNCTIONS

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ABSTRACT. The aim of this article is to describe several cubature formulas related to the Weyl group orbit functions, i.e. to the special cases of the Jacobi polynomials associated to root systems. The diagram containing the relations among the special functions associated to the Weyl group orbit functions is presented and the link between the Weyl group orbit functions and the Jacobi polynomials is explicitly derived in full generality. The four cubature rules corresponding to these polynomials are summarized for all simple Lie algebras and their properties simultaneously tested on model functions. The Clenshaw-Curtis method is used to obtain additional formulas connected with the simple Lie algebra C_2 .

KEYWORDS: Weyl group orbit functions; Jacobi polynomials; cubature formulas.

1. INTRODUCTION

The purpose of this paper is to explicitly overview in full generality the link between the Weyl group orbit functions and the Jacobi and Macdonald polynomials and further examine and compare related methods of numerical integration. These methods of numerical integration, known as cubature rules, emerged recently for all four cases of the Weyl group orbit functions.

The four families of the Weyl group orbit functions [10, 18, 19, 26, 30] are connected to four families of orthogonal polynomials via similar relations between Chebyshev polynomials of the first and second kinds and the ordinary cosine and sine functions. A full set of four families of orbit functions, called C -, S -, S^s - and S^l -functions, arises from root systems of simple Lie algebras with two different lengths of roots. These four families of orthogonal polynomials are in fact special cases of the multivariate Jacobi polynomials [11, 12]. The Jacobi polynomials associated to root systems are in turn limiting cases of the Macdonald polynomials [24]. This connection between the four cases of the Jacobi polynomials and the underlying orbit functions allows to formulate the corresponding methods for numerical integration in terms of the Jacobi polynomials.

Among methods for numerical integration, the quadrature and cubature formulas related to polynomials of a bounded degree hold a prominent place [1, 6–8, 35, 38]. Such formulas estimate a given weighted integral over a fixed domain in Euclidean space. This estimation holds exactly for all polynomials up to a certain degree. A significant effort put into development of various types of cubature formulas results in multitude of types of integration domains with

varying efficiencies. The shapes of the integration domains and the nodes for cubature formulas corresponding to the orthogonal polynomials of the Weyl group orbit functions are determined by the symmetries of the affine Weyl groups and a certain transform [15, 22, 23, 26, 27]. This transform is generated by the transform which induces the given set of orthogonal polynomials. Moreover, a specific notion of the modified degree of multivariate polynomials is essential for establishing the final cubature formulas.

One of the specific methods of deriving quadrature formulas, known as Clenshaw-Curtis method [5], is classically related to Chebyshev polynomials of one variable [9]. Its two-dimensional version related to two-variable Chebyshev polynomials of the root system A_2 is also developed [31]. The importance of this method lies e.g. in its utilization for practical optimization of the shapes of integration domains. The shapes of the integration domains are determined by the underlying Lie algebra [15] and are, however, of non-standard form. In case of simple Lie algebras related to two-variable functions, one of the possible optimizations of these shapes is, similarly to [31], inscribing a triangle into the original fundamental domain. The focus of the present article is on simple Lie algebra C_2 and its corresponding cubature rules.

The integration domain in the case of C_2 is a region bounded by two lines and a parabola depicted in Fig. 4. Except from a general perspective in [15, 26, 27], integration over this region is studied in [32]. Similarly to [31] for A_2 , the non-standard shape of this integration domain motivates further exploration of the Clenshaw-Curtis method. This method crucially depends on the choice of the weight function and the inscribed integration region and has not yet been studied in detail for

the case of C_2 . In this case, the domain inscribed in the original integration region is considered either the original domain itself or the triangle depicted in Fig. 6. Two fundamental choices of the weight functions are detailed. Prior to practical implementation, exact values of certain integrals are also needed. One of the goals of this article is to provide all data necessary for practical implementation of these cubature formulas. This is achieved by tabulating and calculating all needed stabilizer coefficients and exact integral values for each choice of the inscribed integration domain and weight function. To demonstrate usefulness and viability of the presented methods, the numerical tests on model functions, including multidimensional step-functions, are also performed.

The development of novel cubature formulas is motivated by their widespread use in applied numerical simulations and engineering problems. Among direct numerical applications of the cubature formulas is the induced method of polynomial approximation. The cubature formulas are ubiquitous in the modern theory of electromagnetism, especially in its branches of electromagnetic wave propagation [33], magneto-static modeling [39] and micromagnetic simulations [4]. Other fields include fluid flows simulations, laser optics and stochastic dynamics.

In Section 2 are reviewed the notions necessary for definition of the Weyl group orbit functions. The relation between the orbit functions and the Jacobi polynomials is detailed. In Section 3, the cubature formulas from [15, 26, 27] are summarized, Clenshaw-Curtis method is described and used to derive additional cubature formulas. Furthermore, numerical test results are presented.

2. SPECIAL FUNCTIONS ASSOCIATED TO ROOT SYSTEMS

2.1. BASIC DEFINITIONS

This section reviews the basic concepts and notation from the theory of root systems, Weyl groups and Weyl group orbit functions. It is consistent with the notation used in recent papers regarding the topic, such as [10, 13–15, 26] and others. We consider simple Lie algebras, i.e. four infinite families $A_n (n \geq 1)$, $B_n (n \geq 3)$, $C_n (n \geq 2)$ and $D_n (n \geq 4)$ and five exceptional algebras E_6, E_7, E_8, F_4 and G_2 (for the classification see [2, 16]). In particular, we focus on the algebra C_2 as the simplest non-trivial example. Each simple Lie algebra is completely described by its set of simple roots $\Delta = \{\alpha_1, \dots, \alpha_n\}$ which forms a non-orthogonal basis of the Euclidean space \mathbb{R}^n equipped with a scalar product denoted by $\langle \cdot, \cdot \rangle$. Simple roots are either of the same length or of two different lengths, in the latter case we distinguish so-called short and long roots and write $\Delta = \Delta_s \cup \Delta_l$.

The set of dual roots is denoted by

$$\Delta^\vee = \{\alpha_1^\vee, \dots, \alpha_n^\vee\},$$

where $\alpha_i^\vee = \frac{2\alpha_i}{\langle \alpha_i, \alpha_i \rangle}$. In addition to the bases of simple roots and dual roots we introduce the weight basis $\omega_1, \dots, \omega_n$ and the dual weight basis $\omega_1^\vee, \dots, \omega_n^\vee$ where $\langle \alpha_i^\vee, \omega_j \rangle = \langle \alpha_i, \omega_j^\vee \rangle = \delta_{ij}$. The Cartan matrix C is defined as $C_{ij} = \frac{2\langle \alpha_i, \alpha_j \rangle}{\langle \alpha_j, \alpha_j \rangle}$ and its determinant is denoted by c .

Each simple root α_i relates to a reflection r_i defined for every $a \in \mathbb{R}^n$ as

$$r_i a = a - \frac{2\langle a, \alpha_i \rangle}{\langle \alpha_i, \alpha_i \rangle} \alpha_i.$$

The set of reflections $\{r_1, \dots, r_n\}$ generates a finite group W called the Weyl group. By the action of W on the set of simple roots we obtain the root system $\Pi = W\Delta$. Analogously, we define $\Pi^\vee = W\Delta^\vee$, $\Pi^s = W\Delta_s$ and $\Pi^l = W\Delta_l$. Every element of Π can be written as a combination of simple roots with only non-negative (positive roots) or non-positive integer coefficients (negative roots). The set of positive roots is denoted by Π_+ . We define a partial ordering \preceq of roots, $\mu \preceq \lambda$ if $\mu - \lambda$ is a sum of simple roots with non-negative integer coefficients. There is a unique highest root ξ with respect to this ordering, its coordinates in the basis of simple roots are called marks and denoted by m_1, \dots, m_n . Dual root system Π^\vee contains the highest root $\eta = m_1^\vee \alpha_1^\vee + \dots + m_n^\vee \alpha_n^\vee$ with the coefficients m_i^\vee called dual marks.

An infinite extension of the Weyl group W is the affine Weyl group W^{aff} which is obtained by adding to the set of generators of W the affine reflection r_0 ,

$$r_0 a = r_\xi a + \frac{2\xi}{\langle \xi, \xi \rangle}, \quad r_\xi a = a - \frac{2\langle a, \xi \rangle}{\langle \xi, \xi \rangle} \xi.$$

It can also be written as a semidirect product of W and a set of shifts by integer combinations of dual roots [13]. We denote by ψ the retraction homomorphism $W^{\text{aff}} \rightarrow W$ [14]. The fundamental domain F - a set containing exactly one point from each W^{aff} orbit - can be chosen as

$$F = \{ b_1 \omega_1^\vee + \dots + b_n \omega_n^\vee \mid b_i \in \mathbb{R}^{\geq 0}, b_0 + b_1 m_1 + \dots + b_n m_n = 1 \}. \quad (1)$$

Analogously we define dual affine Weyl group as a semidirect product of W and shifts by integer combinations of simple roots.

We introduce three lattices P, P^+ and P^\vee as

$$\begin{aligned} P &= \mathbb{Z}\omega_1 + \dots + \mathbb{Z}\omega_n, \\ P^+ &= \mathbb{Z}^{\geq 0}\omega_1 + \dots + \mathbb{Z}^{\geq 0}\omega_n, \\ P^\vee &= \mathbb{Z}\omega_1^\vee + \dots + \mathbb{Z}\omega_n^\vee. \end{aligned}$$

Note that the root system Π is contained in P , therefore, the partial ordering \preceq can be extended to the lattice P .

A function $k: \alpha \in \Pi \rightarrow k_\alpha \in \mathbb{R}^{\geq 0}$ such that

$$k_\alpha = k_{w(\alpha)} \quad \text{for all } w \in W$$

is known as a multiplicity function on Π . The trivial example is to take $k_\alpha = \text{const}$ for all $\alpha \in \Pi$ which we denote by k^{const} . For simple Lie algebras with two different root lengths, it is natural to distinguish between short and long roots by defining

$$k_\alpha^s \equiv \begin{cases} 1 & \text{if } \alpha \in \Pi^s, \\ 0 & \text{if } \alpha \in \Pi^l, \end{cases} \quad k_\alpha^l \equiv \begin{cases} 0 & \text{if } \alpha \in \Pi^s, \\ 1 & \text{if } \alpha \in \Pi^l. \end{cases}$$

The notion of multiplicity function allows us to define sums of positive roots $\varrho(k)$ and numbers $h(k)$,

$$\varrho(k) \equiv \frac{1}{2} \sum_{\alpha \in \Pi_+} k_\alpha \alpha, \tag{2}$$

$$h(k) = k_\xi + \sum_{i=1}^n m_i k_{\alpha_i}.$$

In particular, with the choice of k^t , where t is one of the symbols $\{0, 1, s, l\}$, we have

$$\begin{aligned} \varrho^0 &\equiv \varrho(k^0) = 0, \\ \varrho^1 &\equiv \varrho(k^1) = \sum_{i=1}^n \omega_i, \\ \varrho^s &\equiv \varrho(k^s) = \sum_{\alpha_i \in \Delta_s} \omega_i, \\ \varrho^l &\equiv \varrho(k^l) = \sum_{\alpha_i \in \Delta_l} \omega_i \end{aligned} \tag{3}$$

and

$$\begin{aligned} h^0 &\equiv h(k^0) = 0, \\ h^1 &\equiv h(k^1) = 1 + \sum_{i=1}^n m_i, \\ h^s &\equiv h(k^s) = \sum_{\alpha_i \in \Delta_s} m_i, \\ h^l &\equiv h(k^l) = 1 + \sum_{\alpha_i \in \Delta_l} m_i. \end{aligned} \tag{4}$$

The number $h \equiv h^1$ is called the Coxeter number, analogously, we call h^s and h^l short and long Coxeter number.

The set of simple roots $\Delta = \{\alpha_1, \alpha_2\}$ of the algebra C_2 decomposes into the set of the short simple roots $\Delta_s = \{\alpha_1\}$ and the set of the long simple roots $\Delta_l = \{\alpha_2\}$. The highest root is of the form $\xi = 2\alpha_1 + \alpha_2$ and the dual highest root is $\eta = \alpha_1^\vee + 2\alpha_2^\vee$. Thus, the sets of mark and dual marks are $(m_1, m_2) = (2, 1)$ and $(m_1^\vee, m_2^\vee) = (1, 2)$. The vectors ϱ^t are of the form $(\varrho^1, \varrho^s, \varrho^l) = (\omega_1 + \omega_2, \omega_1, \omega_2)$ and the Coxeter numbers are $(h^1, h^s, h^l) = (4, 2, 2)$. The roots and dual roots, the weights and dual weights, together with the fundamental domain F and the vectors ϱ^t are depicted in Figure 1.

2.2. WEYL GROUP ORBIT FUNCTIONS

The definition of Weyl group orbit function uses the notion of a sign homomorphisms $\sigma^t : W \mapsto \pm 1$, where

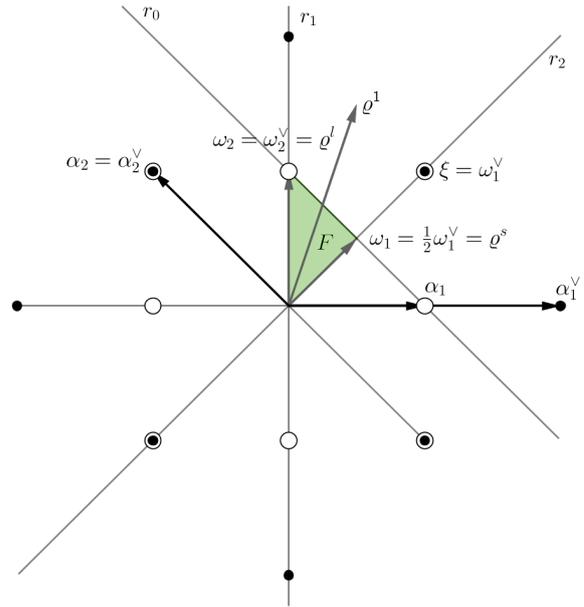


FIGURE 1. Root system of C_2 . The white circles denote the roots, black dots depict the dual roots. The triangle denotes the fundamental domain F . The lines denoted r_1 , r_2 and r_0 depict reflecting mirrors which realize the corresponding reflections.

$t \in \{0, 1, s, l\}$. These can be defined by the values on the reflections corresponding to the simple roots r_i , namely

$$\begin{aligned} \sigma^0(r_i) &= 1, & \alpha_i \in \Delta, \\ \sigma^1(r_i) &= -1, & \alpha_i \in \Delta, \\ \sigma^s(r_i) &= \begin{cases} 1 & \text{if } \alpha_i \in \Delta_l, \\ -1 & \text{if } \alpha_i \in \Delta_s, \end{cases} \\ \sigma^l(r_i) &= \begin{cases} 1 & \text{if } \alpha_i \in \Delta_s, \\ -1 & \text{if } \alpha_i \in \Delta_l. \end{cases} \end{aligned}$$

Several families of special functions are connected with each Weyl group W . They are labelled by vectors $\lambda \in P^+$ and defined as weighed sums over the corresponding Weyl group orbit, i.e. the set $W\lambda = \{w\lambda \mid w \in W\}$. For every $x \in \mathbb{R}^n$ and $t \in \{0, 1, s, l\}$ we define

$$S_{\lambda+\varrho^t}^t(x) = \sum_{\mu \in W(\lambda+\varrho^t)} \sigma^t(\mu) e^{2\pi i \langle \mu, x \rangle}, \tag{5}$$

where ϱ^t is given by (3) and $\sigma^t(\mu) \equiv \sigma^t(w)$ for w such that $\mu = w(\lambda + \varrho^t)$. Functions corresponding to the choice of $t = 0$ and $t = 1$ are usually called C - and S -functions respectively, in the formulas in next sections we use the notation S^0 and S^1 for a simplicity.

Families of C -, S -, S^s - and S^l -functions are complex multivariate functions with remarkable properties such as (anti-)invariance with respect to the action of the affine Weyl group, continuous and discrete orthogonality. They were studied in many papers, see for example [18, 19, 26] for the general properties and [13, 14, 28] for their discretization.

Fundamental domains F^t are defined as subsets of F such that we omit the part of the boundary of F which is stabilized by certain generating reflections $r \in R = \{r_0, r_1, \dots, r_n\}$. More precisely, with the notation

$$R^t = \{r \in R \mid \sigma^t \circ \psi(r) = -1\},$$

$$H^t = \{a \in F \mid \exists r \in R^t, ra = a\},$$

we define $F^t = F \setminus H^t$. The explicit forms are obtained from (1) and can be found in [14].

The C -, S -, S^s - and S^l -functions can be viewed as functional forms of elements from the algebra $\mathbb{C}[P]$ containing all complex linear combinations of formal exponentials $e^a, a \in P$, with multiplication defined by $e^a \cdot e^b = e^{a+b}$, the inverse given by $(e^a)^{-1} = e^{-a}$ and the identity $e^0 = 1$. The connection is based on the exponential mapping from Lie algebra to the corresponding Lie group [2, 16, 27].

2.3. JACOBI POLYNOMIALS

We assume that the multiplicity function k satisfies $k_\alpha \geq 0$. The Jacobi polynomial $P(\lambda, k)$ [11, 12] associated to the root system Π with highest weight $\lambda \in P^+$ and multiplicity function k as parameter is defined by the following formulas.

$$P(\lambda, k) \equiv \sum_{\substack{\mu \in P^+ \\ \mu \preceq \lambda}} c_{\lambda\mu}(k) C_\mu, \quad C_\mu = \sum_{\mu' \in W\mu} e^{\mu'}, \quad (6)$$

where the coefficients $c_{\lambda\mu}(k)$ are recursively given by

$$\begin{aligned} & (\langle \lambda + \varrho(k), \lambda + \varrho(k) \rangle - \langle \mu + \varrho(k), \mu + \varrho(k) \rangle) c_{\lambda\mu}(k) \\ &= 2 \sum_{\alpha \in \Pi_+} k_\alpha \sum_{j=1}^{\infty} \langle \mu + j\alpha, \alpha \rangle c_{\lambda, \mu+j\alpha}(k) \end{aligned}$$

along with the initial value $c_{\lambda\lambda} = 1$ and the assumption $c_{\lambda\mu} = c_{\lambda, w(\mu)}$ for all $w \in W$. Recall that $\varrho(k)$ is defined by (2).

By setting $k_\alpha = 0$ for all $\alpha \in \Pi$, the Jacobi polynomials lead trivially to C -functions. In the case $k = k^1$, the formula for the calculation of the coefficients becomes the Freudenthal's recurrence formula [16]. Therefore, each $P(\lambda, k^1)$ specializes to a character χ_λ of an irreducible representation of the simple Lie algebra of the highest weight λ , i.e.,

$$P(\lambda, k^1) = \chi_\lambda = \frac{S_{\lambda+\varrho^1}}{S_{\varrho^1}}.$$

In addition, we show that the Jacobi polynomials are related to the S^s - and S^l -functions in the following way.

$$P(\lambda, k^s) = \frac{S_{\lambda+\varrho^s}^s}{S_{\varrho^s}^s} \quad \text{and} \quad P(\lambda, k^l) = \frac{S_{\lambda+\varrho^l}^l}{S_{\varrho^l}^l}.$$

We first observe that $S_{\lambda+\varrho^s}^s/S_{\varrho^s}^s$ are Weyl group invariant elements of $\mathbb{C}[P]$ (see Proposition 4.2 of [26])

with well-known basis formed by C -functions [2]. Therefore, each $S_{\lambda+\varrho^s}^s/S_{\varrho^s}^s$ can be expressed as a linear combination of C -functions. By the definition of the Weyl group, the weights of exponentials in $S_{\lambda+\varrho^s}^s$ are of the form $\lambda + \varrho^s - g(\alpha_1, \dots, \alpha_n)$, where $g(\alpha_1, \dots, \alpha_n)$ denotes a sum of simple roots with non-negative integer coefficients, and the unique maximal weight is $\lambda + \varrho^s$. Similarly, the unique maximal weight of $S_{\varrho^s}^s$ is ϱ^s . Therefore, we have

$$\frac{S_{\lambda+\varrho^s}^s}{S_{\varrho^s}^s} = \sum_{\substack{\mu \in P^+ \\ \mu \preceq \lambda}} b_\mu C_\mu, \quad b_\lambda = 1.$$

To prove $b_\mu = c_{\lambda\mu}(k^s)$, we proceed by using an equivalent definition of Jacobi polynomials with the multiplicity function satisfying $k_\alpha \in \mathbb{Z}^{\geq 0}$ [12]. For any $f = \sum_\lambda a_\lambda e^\lambda$, we define

$$\bar{f} \equiv \sum_\lambda \bar{a}_\lambda e^{-\lambda} \quad \text{and} \quad CT(f) \equiv a_0.$$

If we introduce the scalar product (\cdot, \cdot) on $\mathbb{C}[P]$ by

$$\begin{aligned} (f, g) &\equiv CT(f\bar{g}\delta(k)^{\frac{1}{2}}\overline{\delta(k)^{\frac{1}{2}}}), \quad f, g \in \mathbb{C}[P], \\ \delta(k)^{\frac{1}{2}} &\equiv \prod_{\alpha \in \Pi_+} (e^{\frac{1}{2}\alpha} - e^{-\frac{1}{2}\alpha})^{k_\alpha}, \end{aligned}$$

then the Jacobi polynomials $P(\lambda, k)$ are the unique polynomials of the form (6) satisfying the requirement

$$(P(\lambda, k), P(\mu, k)) = 0$$

for all $\mu \in P^+$ such that $\mu \preceq \lambda$ and $\lambda \neq \mu$ assuming $c_{\lambda\lambda} = 1$.

Using Proposition 4.1 of [26], i.e. $\delta(k^s)^{\frac{1}{2}} = S_{\varrho^s}^s$, we obtain

$$\begin{aligned} & \left(\frac{S_{\lambda+\varrho^s}^s}{S_{\varrho^s}^s}, \frac{S_{\mu+\varrho^s}^s}{S_{\varrho^s}^s} \right) = CT(S_{\lambda+\varrho^s}^s \overline{S_{\mu+\varrho^s}^s}) \\ &= CT \left(\sum_{\lambda' \in W(\lambda+\varrho^s)} \sum_{\mu' \in W(\mu+\varrho^s)} \sigma^s(\lambda') \sigma^s(\mu') e^{\lambda' - \mu'} \right). \end{aligned}$$

Clearly $\lambda' = \mu'$ if and only if there exists $w \in W$ such that $\lambda + \varrho^s = w(\mu + \varrho^s)$. For we consider $\lambda \in P^+$ different from $\mu \in P^+$, it is not possible to have $\lambda' = \mu'$. This implies that

$$\left(\frac{S_{\lambda+\varrho^s}^s}{S_{\varrho^s}^s}, \frac{S_{\mu+\varrho^s}^s}{S_{\varrho^s}^s} \right) = 0 \quad \text{and} \quad P(\lambda, k^s) = \frac{S_{\lambda+\varrho^s}^s}{S_{\varrho^s}^s}.$$

The proof for the long root case is similar.

Finally, note that the Jacobi polynomials can be viewed as the limiting case of the Macdonald polynomials $P_\lambda(q, t_\alpha)$ when $t_\alpha = q^{k_\alpha}$ with k_α fixed and $q \rightarrow 1$. See [24] for more details. The relations among several special functions associated with the Weyl groups, which are summarized in [29], are depicted in Figure 2.

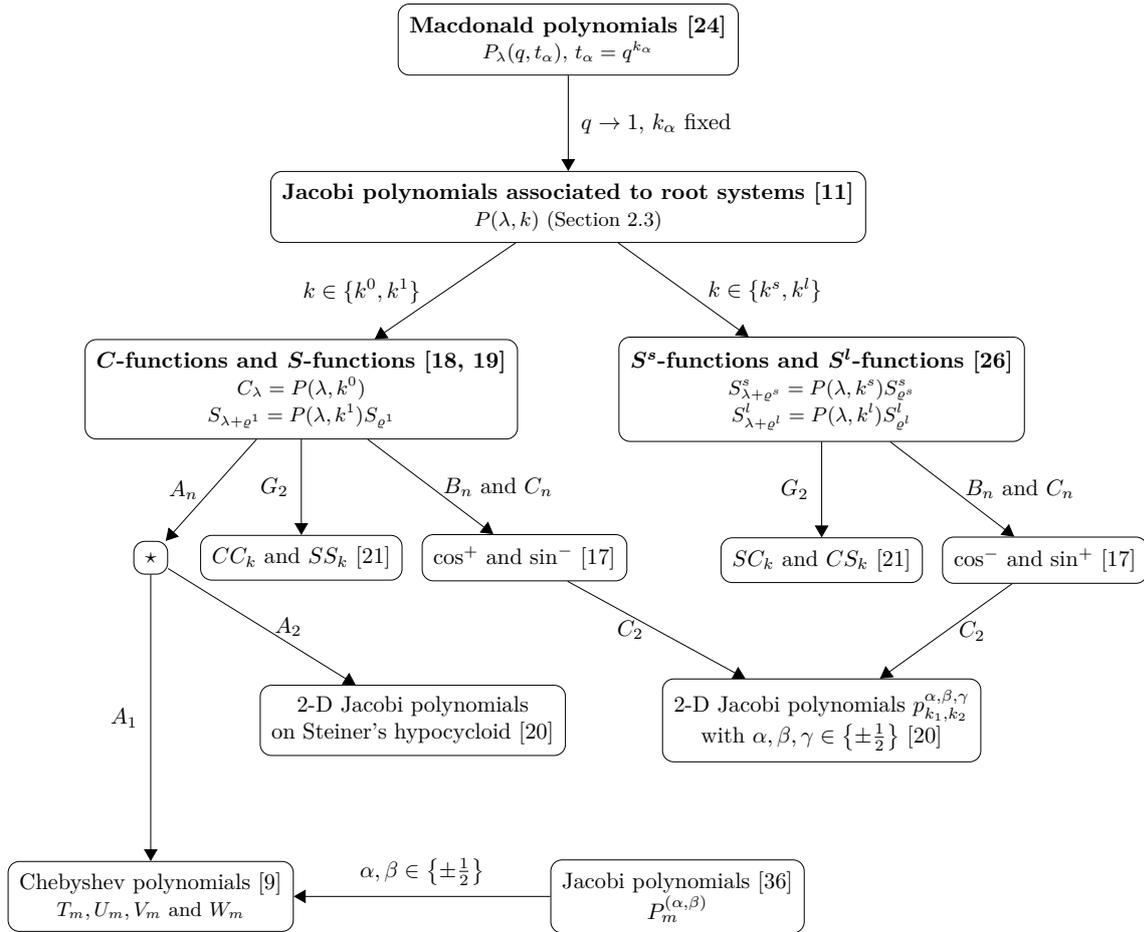


FIGURE 2. The diagram of relations among several special functions associated with Weyl groups.

3. CUBATURE FORMULAS

3.1. GENERAL FORM OF CUBATURE FORMULAS

Analogously to Chebyshev polynomials, we identify polynomial variables y_1, \dots, y_n with real-valued functions in the following way. Let $z_j \equiv C_{\omega_j}$, then

$$\begin{aligned}
 A_{2k}: y_j &= \Re(z_j), y_{2k-j+1} = \Im(z_j), j = 1, \dots, k, \\
 A_{2k+1}: y_j &= \Re(z_j), y_{k+1} = z_{k+1}, y_{2k-j+2} = \Im(z_j), \\
 & j = 1, \dots, k, \\
 D_{2k+1}: y_j &= z_j, y_{2k} = \Re(z_{2k}), \\
 & y_{2k+1} = \Im(z_{2k}), j = 1, \dots, 2k - 1, \\
 E_6: y_1 &= \Re(z_1), y_2 = \Re(z_2), y_3 = z_3, \\
 & y_4 = \Im(z_2), y_5 = \Im(z_1), y_6 = z_6,
 \end{aligned}$$

otherwise we put $y_j = z_j$. We say that a monomial $y_1^{\lambda_1} \dots y_n^{\lambda_n}$ has an m -degree

$$\deg_m y_1^{\lambda_1} \dots y_n^{\lambda_n} = m_1^\vee \lambda_1 + \dots + m_n^\vee \lambda_n$$

and any polynomial p in $\mathbb{C}[y_1, \dots, y_n]$ has an m -degree equal to the largest m -degree of the monomials occurring in p . We denote the subspace containing all polynomials of m -degree at most M by Π_M . For $\lambda = \lambda_1 \omega_1 + \dots + \lambda_n \omega_n$, the C -functions C_λ and thus all Jacobi polynomials $P(\lambda, k)$ can be rewritten as

orthogonal polynomials in variables y_1, \dots, y_n of m -degree equal to $m_1^\vee \lambda_1 + \dots + m_n^\vee \lambda_n$ [15].

The variables y_1, \dots, y_n viewed as functions induce a map

$$\Xi: \mathbb{R}^n \rightarrow \mathbb{R}^n, \quad \Xi(x) = (y_1(x), \dots, y_n(x)).$$

The map Ξ is used to define the integration region Ω and the sets of nodes $\Omega_M^t, t \in \{0, 1, s, l\}$ by

$$\Omega \equiv \Xi(F^1), \quad \Omega_M^t \equiv \Xi\left(\frac{1}{M+h^t} P^\vee \cap F^t\right).$$

Let

$$\begin{aligned}
 \text{Stab}_{W^{\text{aff}}}(x) &\equiv \{w \in W^{\text{aff}} \mid wx = x\}, \\
 \varepsilon(x) &\equiv \frac{|W|}{|\text{Stab}_{W^{\text{aff}}}(x)|},
 \end{aligned}$$

and define a map $\tilde{\varepsilon}: \Omega_M^0 \rightarrow \mathbb{N}$ by

$$\tilde{\varepsilon}(y) \equiv \varepsilon(\Xi_M^{-1}y), \quad \Xi_M = \Xi \upharpoonright_{\frac{1}{M+h^t} P^\vee \cap F^0}.$$

Denoting $K(y_1, \dots, y_n) \equiv \sqrt{S_{\rho^1} \overline{S_{\rho^1}}}$, the weight functions are given by

$$\begin{aligned}
 s^t(y_1, \dots, y_n) &\equiv S_{\rho^t}^t \overline{S_{\rho^t}^t}, \\
 w^t(y_1, \dots, y_n) &\equiv \frac{s^t(y_1, \dots, y_n)}{K(y_1, \dots, y_n)}, \quad t \in \{0, 1, s, l\}.
 \end{aligned}$$

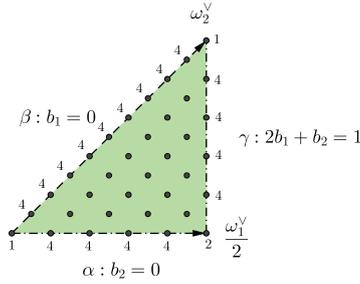


FIGURE 3. The fundamental domain $F = F^0$ of C_2 is depicted as the triangle with dashed boundary H^s and dot-and-dashed boundary H^l . The black dots correspond to the points from $\frac{1}{10}P^\vee \cap F$. The numbers 1, 2, 4 of the dots are the values of $\varepsilon(x)$, the inner dots have $\varepsilon(x) = 8$.

Since the products $S_{\theta^t}^t \overline{S_{\theta^t}^t}$, $t \in \{0, 1, s, l\}$ are W -invariant sums of exponentials from $\mathbb{C}[P]$, they are expressible as functions in polynomial variables y_1, \dots, y_n .

In [15, 26, 27] is shown that the following cubature formulas are exact equalities for any $M \in \mathbb{N}$ and any polynomial p which satisfies the following constraints,

- $\deg_m p \leq 2M - 1$ for $t \in \{0, l\}$,
- $\deg_m p \leq 2M + 1$ for $t \in \{1, s\}$.

Thus, it holds that

$$\int_{\Omega} p(y)w^t(y) dy = \frac{\kappa}{c|W|} \left(\frac{2\pi}{M+h^t} \right)^n \sum_{y \in \Omega_M^t} \tilde{\varepsilon}(y) s^t(y) p(y), \quad (7)$$

where

$$\kappa = \begin{cases} 2^{-\lfloor \frac{n}{2} \rfloor} & \text{for } A_n, \\ \frac{1}{2} & \text{for } D_{2k+1}, \\ \frac{1}{4} & \text{for } E_6, \\ 1 & \text{otherwise.} \end{cases}$$

To numerically compare the efficiency of these cubature formulas, we may consider an integrable function f , such that f/w^t is well defined on Ω_M^t , and rewrite the cubatures (7) in the following form:

$$I(f) \equiv \int_{\Omega} f(y) dy \approx I_M^t(f), \quad (8)$$

$$I_M^t(f) \equiv \frac{\kappa}{c|W|} \left(\frac{2\pi}{M+h^t} \right)^n \sum_{y \in \Omega_M^t} \tilde{\varepsilon}(y) K(y) f(y).$$

3.2. CUBATURE FORMULAS OF C_2

In this section, the general cubature formulas (8) are specialized and tested on model examples for the case of algebra C_2 . The region F^t of C_2 with the points from $\frac{1}{10}P^\vee \cap F^t$ is depicted in Fig. 3, whereas the corresponding integration region Ω with the transformed grid points is depicted in Fig. 4. Note that

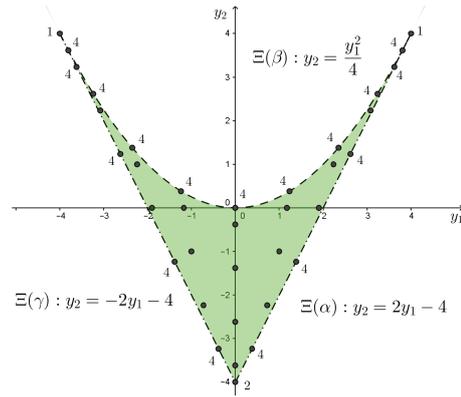


FIGURE 4. The integration region Ω of C_2 contains the points of the grid Ω_{10}^0 . The inner points of Ω corresponds to the grid Ω_6^1 , the points not lying on the dashed boundary corresponds to the grid Ω_8^s and finally, the points not lying on the dot-and-dashed boundary corresponds to the grid Ω_8^l . The numbers 1, 2, 4 are the values of $\tilde{\varepsilon}(y)$, the inner dots have $\tilde{\varepsilon}(y) = 8$.

the numbers of points in grids Ω_M^t , $t \in \{0, 1, s, l\}$ are the same. Fixing the basis $x = b_1\omega_1^\vee + b_2\omega_2^\vee$ results in the polynomial variables expressed as

$$y_1 = 2(\cos \pi(2b_1 + b_2) + \cos \pi b_2),$$

$$y_2 = 2(\cos 2\pi(b_1 + b_2) + \cos 2\pi b_1).$$

Formula (8) specializes into

$$I_M^t(f) = \frac{\pi^2}{4(M+h^t)^2} \sum_{y \in \Omega_M^t} \tilde{\varepsilon}(y) K(y) f(y), \quad (9)$$

where:

- $M \in \mathbb{N}$ is arbitrary;
- $h^0 = 1, h^1 = 4$ and $h^s = h^l = 2$;
- the integration region Ω , depicted in Fig. 4, is bounded by two lines $y_2 = \pm y_1 - 4$ and the parabola $y_2 = \frac{y_1^2}{4}$;
- the finite grid Ω_M^t , depicted for $M = 10$ in Fig. 4), consists of points $(y_1(x), y_2(x))$ where $x = \frac{s_1^t}{M+h^t}\omega_1^\vee + \frac{s_2^t}{M+h^t}\omega_2^\vee$ with s_i^t satisfying

$$s_i^0 \in \mathbb{Z}^{\geq 0}, \quad 2s_1^0 + s_2^0 \leq M,$$

$$s_i^1 \in \mathbb{Z}^{> 0}, \quad 2s_1^1 + s_2^1 < M + 4,$$

$$s_1^s \in \mathbb{Z}^{\geq 0}, s_2^s \in \mathbb{Z}^{> 0}, \quad 2s_1^s + s_2^s \leq M + 2,$$

$$s_1^l \in \mathbb{Z}^{> 0}, s_2^l \in \mathbb{Z}^{\geq 0}, \quad 2s_1^l + s_2^l < M + 2;$$

- the weight function K becomes

$$K(y_1, y_2) = \sqrt{(y_1^2 - 4y_2)((y_2 + 4)^2 - 4y_1^2)};$$

- the weight function $\tilde{\varepsilon}$ is equal to

$$\tilde{\varepsilon}(y) = \begin{cases} 1 & \text{if } (y_1, y_2) = (\pm 4, 4), \\ 2 & \text{if } (y_1, y_2) = (0, -4), \\ 8 & \text{if } (y_1, y_2) \text{ is an inner point of } \Omega, \\ 4 & \text{otherwise.} \end{cases}$$

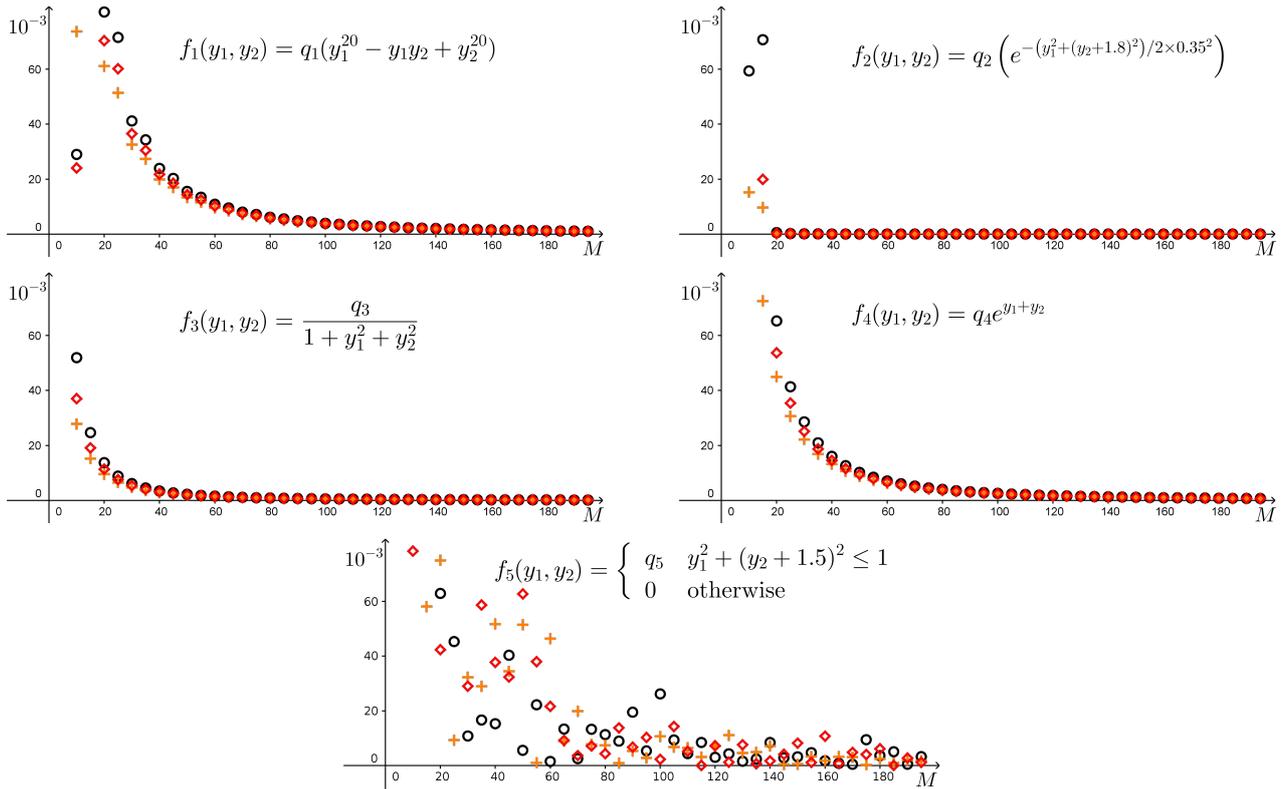


FIGURE 5. The graphs of error values $|1 - I_M^t(f_i)|$ of the integral $I(f_i) = 1$ and its estimations $I_M^t(f_i)$, $M = 10, 15, 20, \dots, 195$ given by (8). The values for $t = 0, 1, s$ are depicted as circles, “+” signs and diamonds, respectively.

For the purpose of numerical tests and comparison, we choose

$$\begin{aligned} f_1(y_1, y_2) &= q_1(y_1^{20} - y_1y_2 + y_2^{20}), \\ f_2(y_1, y_2) &= q_2 \left(e^{-(y_1^2 + (y_2 + 1.8)^2) / 2 \times 0.35^2} \right), \\ f_3(y_1, y_2) &= \frac{q_3}{1 + y_1^2 + y_2^2}, \\ f_4(y_1, y_2) &= q_4 e^{y_1 + y_2}, \\ f_5(y_1, y_2) &= \begin{cases} q_5 & \text{if } y_1^2 + (y_2 + 1.5)^2 \leq 1, \\ 0 & \text{otherwise.} \end{cases} \end{aligned}$$

as model functions. Each value of $q_i \in \mathbb{R}$ is set to satisfy the condition $I(f_i) = 1$. Fig. 5 shows for $M = 10, 15, 20, \dots, 195$ and $t \in \{0, 1, s\}$ the graphs of the absolute value of the difference $|1 - I_M^t(f_i)|$. Note that the cases $t = s$ and $t = l$ give the same results since $h^s = h^l = 2$ and $K(y)f_i(y)$ vanish on the boundary of Ω .

4. CLENSHAW-CURTIS CUBATURE FORMULAS

4.1. CLENSHAW-CURTIS METHOD

Assuming that we have an interpolation of a function f in terms of $P(\lambda, k^t) \in \Pi_M$ in points Ω_M^t , i.e.

$$f \approx \sum_{\substack{\lambda \in P^+ \\ \langle \lambda, \eta \rangle \leq M}} b_\lambda^t P(\lambda, k^t),$$

$$f(y) = \sum_{\substack{\lambda \in P^+ \\ \langle \lambda, \eta \rangle \leq M}} b_\lambda^t P(\lambda, k^t; y), \quad y \in \Omega_M^t,$$

we estimate a weighted integral of f with a weight function w over a domain $D \subset \Omega$ by

$$\sum_{\substack{\lambda \in P^+ \\ \langle \lambda, \eta \rangle \leq M}} b_\lambda^t \int_D P(\lambda, k^t; y) w(y) dy.$$

Such construction of the Clenshaw-Curtis cubature rule implies the exact equality for any polynomial f of m -degree at most M . Denoting

$$a_\lambda^t(w) \equiv \int_D P(\lambda, k^t; y) w(y) dy,$$

the Clenshaw-Curtis cubature is thus given by

$$\int_D f(y) w^t(y) dy \approx \sum_{\substack{\lambda \in P^+ \\ \langle \lambda, \eta \rangle \leq M}} b_\lambda^t a_\lambda^t(w),$$

where the coefficients b_λ^t and $a_\lambda^t(w)$ need to be determined. The coefficients b_λ^t are readily obtained using the discrete orthogonality relations of the orbit functions from [13, 14]. Denoting the order of the stabilizer of $\frac{\lambda + \varrho^t}{M + h^t}$ with respect to the dual affine Weyl group by $h_{\lambda + \varrho^t}^\vee$, it holds that

$$\begin{aligned} b_\lambda^t &= \frac{|\text{Stab}_W(\lambda + \varrho^t)|^2}{c|W|(M + h^t)^n h_{\lambda + \varrho^t}^\vee} \\ &\quad \times \sum_{y \in \Omega_M^t} \tilde{\varepsilon}(y) s^t(y) \overline{f(y) P(\lambda, k^t; y)}. \end{aligned} \quad (10)$$

| $A_{(\lambda_1, \lambda_2)}^0$ | |
|---|---|
| $(\lambda_1, \lambda_2) = (2i, 2j)$ | $\sum_{(\mu_1, \mu_2) \in \mathcal{M}(2i, 2j)} \frac{64(4\mu_2^2 + 4\mu_1\mu_2 - 3)}{ \text{Stab}_W(2i, 2j) (\mu_1^2 - 1)(4\mu_2^2 - 1)(4\mu_2^2 - 9)}$ |
| $(\lambda_1, \lambda_2) = (2i, 2j + 1)$ | $\sum_{(\mu_1, \mu_2) \in \mathcal{M}(2i, 2j+1)} \frac{-64(4\mu_2^2 + 4\mu_1\mu_2 + 3)}{ \text{Stab}_W(2i, 2j+1) (\mu_1^2 - 4)(4\mu_2^2 - 1)(4\mu_2^2 - 9)}$ |
| Otherwise | 0 |
| $A_{(\lambda_1, \lambda_2)}^1$ | |
| $(\lambda_1, \lambda_2) = (2i, 2j)$ | $\frac{32(i+j+1)}{(2i+4j+3)(2j+1)(2i+1)}$ |
| $(\lambda_1, \lambda_2) = (2i, 2j + 1)$ | $\frac{32(j+1)}{(2i+4j+5)(2i+2j+3)(2i+1)}$ |
| Otherwise | 0 |
| $A_{(\lambda_1, \lambda_2)}^s$ | |
| $(\lambda_1, \lambda_2) = (2i, 2j)$ | $\sum_{(\mu_1, \mu_2) \in \mathcal{M}_1^s(2i+1, 2j)} \frac{8(\mu_1 + \mu_2)}{ \text{Stab}_W(2i+1, 2j) \mu_2(\mu_1^2 - 1)(\mu_2^2 - 1)}$ |
| $(\lambda_1, \lambda_2) = (2i, 2j + 1)$ | $\sum_{(\mu_1, \mu_2) \in \mathcal{M}_2^s(2i+1, 2j+1)} \frac{-8(\mu_1 + \mu_2)}{\mu_2(\mu_1^2 - 1)(\mu_2^2 - 1)}$ |
| Otherwise | 0 |
| $A_{(\lambda_1, \lambda_2)}^l$ | |
| $(\lambda_1, \lambda_2) = (2i, 2j)$ | $\left(\sum_{(\mu_1, \mu_2) \in \mathcal{M}_1^l(2i, 2j+1)} - \sum_{(\mu_1, \mu_2) \in \mathcal{M}_2^l(2i, 2j+1)} \right) \frac{32\mu_2}{ \text{Stab}_W(2i, 2j+1) \mu_1(4\mu_2^2 - 1)}$ |
| $(\lambda_1, \lambda_2) = (2i, 2j + 1)$ | $\left(\sum_{(\mu_1, \mu_2) \in \mathcal{M}_2^l(2i, 2j+2)} - \sum_{(\mu_1, \mu_2) \in \mathcal{M}_1^l(2i, 2j+2)} \right) \frac{16(2\mu_1\mu_2 + 1)}{ \text{Stab}_W(2i, 2j+2) (\mu_1^2 - 1)(4\mu_2^2 - 1)}$ |
| Otherwise | 0 |
| $\mathcal{M}(\lambda_1, \lambda_2)$ | $\{(\lambda_1 + \lambda_2, \frac{\lambda_1}{2} + \lambda_2), (\lambda_2, \frac{\lambda_1}{2} + \lambda_2), (\lambda_1 + \lambda_2, \frac{\lambda_1}{2}), (\lambda_2, -\frac{\lambda_1}{2})\}$ |
| $\mathcal{M}_1^s(\lambda_1, \lambda_2)$ | $\{(\lambda_2, \frac{\lambda_1}{2} + \lambda_2), (\lambda_2, -\frac{\lambda_1}{2})\}$ |
| $\mathcal{M}_2^s(\lambda_1, \lambda_2)$ | $\{(\lambda_1 + \lambda_2, \frac{\lambda_1}{2} + \lambda_2), (\lambda_1 + \lambda_2, \frac{\lambda_1}{2})\}$ |
| $\mathcal{M}_1^l(\lambda_1, \lambda_2)$ | $\{(\lambda_1 + \lambda_2, \frac{\lambda_1}{2} + \lambda_2), (\lambda_2, \frac{\lambda_1}{2} + \lambda_2)\}$ |
| $\mathcal{M}_2^l(\lambda_1, \lambda_2)$ | $\{(\lambda_1 + \lambda_2, \frac{\lambda_1}{2}), (\lambda_2, -\frac{\lambda_1}{2})\}$ |

TABLE 2. Values of A_λ^t (11) for $t \in \{0, 1, s, l\}$, $\lambda = \lambda_1\omega_1 + \lambda_2\omega_2$ and i, j are non-negative integers.

| $[\lambda_0, \lambda_1, \lambda_2]$ | $ \text{Stab}_W(\lambda + \varrho^t) $ | $h_{\lambda + \varrho^t}^\vee$ |
|-------------------------------------|--|--------------------------------|
| (\star, \star, \star) | 1 | 1 |
| $(0, \star, \star)$ | 1 | 2 |
| $(\star, 0, \star)$ | 2 | 2 |
| $(\star, \star, 0)$ | 2 | 2 |
| $(0, 0, \star)$ | 2 | 4 |
| $(0, \star, 0)$ | 2 | 8 |
| $(\star, 0, 0)$ | 8 | 8 |

TABLE 1. The values of $|\text{Stab}_W(\lambda + \varrho^t)|$ and $h_{\lambda + \varrho^t}^\vee$ of C_2 , where $\lambda + \varrho^t = \lambda_1\omega_1 + \lambda_2\omega_2$ and $\lambda_0 \equiv M + h^t - \lambda_1 - 2\lambda_2$. Asterisks denote non-zero positive integers.

It remains to evaluate the integrals $a_\lambda^t(w)$ which depend on the chosen weight and the integration domain $D \subset \Omega$.

Since the Jacobi polynomials have several properties connected to the domain Ω (e.g. continuous and discrete orthogonality), we firstly take $D = \Omega$. The cubature rules with the choice $w = w^t$ coincide for any simple Lie algebra with the formulas (7). The difference lies in the fact that Clenshaw-Curtis method guarantees the exact equality only for polynomials up to m -degree M .

4.2. INTEGRATION DOMAIN Ω OF C_2

In this section, the Clenshaw-Curtis integration method is applied to the algebra C_2 . The values of $|\text{Stab}_W(\lambda + \varrho^t)|$ and $h_{\lambda + \varrho^t}^\vee$, needed in (10), are tabulated in Tab. 1.

Since the choice of $w = w^t$ gives standard cubature formulas, the next natural choice of the weight function is to set $w = 1$. In this case are the coefficients $a_\lambda^t(1)$, denoted by A_λ^t , expressed as the following integrals:

$$A_\lambda^t = 2\pi^2 \begin{cases} \int_F C_\lambda(x) S_{\varrho^1}(x) dx & \text{if } t = 0, \\ \int_F S_{\lambda + \varrho^1}(x) dx & \text{if } t = 1, \\ \int_F S_{\lambda + \varrho^s}^s(x) S_{\varrho^l}^l(x) dx & \text{if } t = s, \\ \int_F S_{\lambda + \varrho^l}^l(x) S_{\varrho^s}^s(x) dx & \text{if } t = l. \end{cases} \quad (11)$$

The exact values of A_λ^t are explicitly calculated in Tab. 2.

4.3. TRIANGULAR DOMAIN OF C_2

The next choice of the domain D , for which we derive the Clenshaw-Curtis cubature rules, is the triangle $T \subset \Omega$ depicted on Fig. 6 and given explicitly by

$$T \equiv \left\{ (y_1, y_2) \mid y_2 \leq 0, -\frac{y_2}{2} - 2 \leq y_1 \leq \frac{y_2}{2} + 2 \right\}.$$

| $\mathcal{A}_{(\lambda_1, \lambda_2)}^0$ | |
|---|--|
| $(\lambda_1, \lambda_2) = (0, 0)$ | $\frac{\pi^2}{4}$ |
| $(\lambda_1, \lambda_2) = (2i, 2j + 1)$ | $(-1)^{i+1} \frac{8}{ \text{Stab}_W(2i, 2j+1) (2i+2j+1)(2j+1)}$ |
| Otherwise | 0 |
| $\mathcal{A}_{(\lambda_1, \lambda_2)}^1$ | |
| $(\lambda_1, \lambda_2) = (0, 0)$ | $2\pi^2 + \frac{128}{9}$ |
| $(\lambda_1, \lambda_2) = (2i, 2j), i + j \neq 0$ | $(-1)^{i+1} \frac{128(i+j+1)}{(2i+2j+1)(2j-1)(2i+2j+3)(2j+3)}$ |
| $(\lambda_1, \lambda_2) = (2i, 2j + 1)$ | $(-1)^{i+1} \frac{128(j+1)}{(2j+1)(2i+2j+1)(2j+3)(2i+2j+5)}$ |
| Otherwise | 0 |
| $\mathcal{A}_{(\lambda_1, \lambda_2)}^s$ | |
| $(\lambda_1, \lambda_2) = (0, 0)$ | $\pi^2 + 8$ |
| $(\lambda_1, \lambda_2) = (2i, 2j), i + j \neq 0$ | $(-1)^{i+1} \frac{16}{ \text{Stab}_W(2i+1, 2j) (2j+1)(2i+2j+1)(2j-1)}$ |
| $(\lambda_1, \lambda_2) = (2i, 2j + 1)$ | $(-1)^{i+1} \frac{16}{(2j+1)(2i+2j+3)(2i+2j+1)}$ |
| Otherwise | 0 |
| $\mathcal{A}_{(\lambda_1, \lambda_2)}^l$ | |
| $(\lambda_1, \lambda_2) = (0, 0)$ | π^2 |
| $(\lambda_1, \lambda_2) = (2i, 2j + 1)$ | $(-1)^{i+1} \frac{128(i+j+1)(j+1)}{ \text{Stab}_W(2i, 2j+2) (2i+2j+3)(2j+3)(2i+2j+1)(2j+1)}$ |
| Otherwise | 0 |

TABLE 3. Values of \mathcal{A}_λ^t , given by (12), for $t \in \{0, 1, s, l\}$ and $\lambda = \lambda_1\omega_1 + \lambda_2\omega_2$. The indices i, j are non-negative integers.

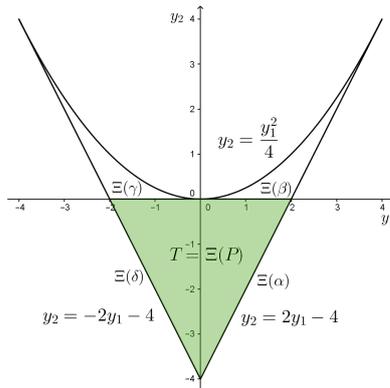


FIGURE 6. The domain bounded by the two lines and the parabola with the inscribed triangle T corresponds to the integration region Ω of C_2 .

Choosing the weight function $w = w^t$, the integrals $a_\lambda^t(w^t)$ are calculated by a change of variables induced by the map Ξ . Denoting $\mathcal{A}_\lambda^t \equiv a_\lambda^t(w^t)$, it holds that

$$\mathcal{A}_\lambda^t = 2\pi^2 \int_P S_{\lambda+\varrho^t}^t(x) \overline{S_{\varrho^t}^t(x)} dx, \quad (12)$$

where P is the pre-image of the triangle T under the map Ξ . This pre-image P , depicted as a square in Fig. 7, contains the points $b_1\omega_1^\vee + b_2\omega_2^\vee$ satisfying $2b_1 + b_2 \geq 1/2$, $2b_1 + b_2 \leq 1$, $b_2 \geq 0$ and $b_2 \leq 1/2$. The exact values of \mathcal{A}_λ^t are tabulated in Tab. 3.

Finally, choosing $w = 1$, we calculate the coefficients

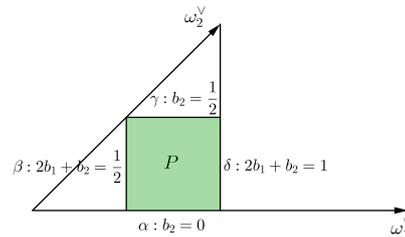


FIGURE 7. The fundamental domain F corresponding to C_2 is depicted as the triangle containing the square P with the boundaries α, β, γ and δ .

$\mathcal{B}_\lambda^t \equiv a_\lambda^t(1)$ on T as the following integrals

$$\mathcal{B}_\lambda^t = 2\pi^2 \begin{cases} \int_P C_\lambda(x) S_{\varrho^1}(x) dx & \text{if } t = 0, \\ \int_P S_{\lambda+\varrho^1}(x) dx & \text{if } t = 1, \\ \int_P S_{\lambda+\varrho^s}^s(x) S_{\varrho^t}^l(x) dx & \text{if } t = s, \\ \int_P S_{\lambda+\varrho^l}^l(x) S_{\varrho^s}^s(x) dx & \text{if } t = l. \end{cases} \quad (13)$$

The exact values of \mathcal{B}_λ^t are tabulated in Tab. 4.

We choose the following functions as model functions for numerical tests,

$$\begin{aligned} g_1(y_1, y_2) &= r_1(y_1^{20} - y_1y_2 + y_2^{20}), \\ g_2(y_1, y_2) &= r_2(e^{-(y_1^2+(y_2+1.8)^2)/2 \times 0.35^2}), \\ g_3(y_1, y_2) &= \frac{r_3}{1 + y_1^2 + y_2^2}, \\ g_4(y_1, y_2) &= r_4e^{x+y}, \end{aligned}$$

| $\mathcal{B}_{(\lambda_1, \lambda_2)}^0$ | |
|---|--|
| $(\lambda_1, \lambda_2) = (0, 1)$ | $-\frac{32}{3}$ |
| $(\lambda_1, \lambda_2) = (2, 0)$ | $-\frac{16}{3}$ |
| $(\lambda_1, \lambda_2) = (2i, 2), i \neq 0$ | $\frac{16(1+(-1)^{i+1})}{3i(i+1)}$ |
| $(\lambda_1, \lambda_2) = (2i, 1), i \neq 0$ | $16 \left[\frac{2}{(2i+3)(2i-1)} + \frac{1+(2i+1)(-1)^{i+1}}{3i(i+1)} \right]$ |
| $(\lambda_1, \lambda_2) = (2i, 2j+1), j \neq 0$ | $\frac{64}{ \text{Stab}_W(2i, 2j+1) } \left[\frac{-1+(2j+1)(-1)^j}{((2i+2j+1)^2-4)((2j+1)^2-1)} + \frac{-1+(2i+2j+1)(-1)^{i+j}}{((2i+2j+1)^2-1)((2j+1)^2-4)} \right]$ |
| $(\lambda_1, \lambda_2) = (2i, 2j), j \neq 1, i+j \neq 1$ | $\frac{16}{ \text{Stab}_W(2i, 2j) } \left[\frac{1+(-1)^{i+j}}{((i+j)^2-1)(4j^2-1)} + \frac{1+(-1)^j}{(4(i+j)^2-1)(j^2-1)} \right]$ |
| Otherwise | 0 |
| $\mathcal{B}_{(\lambda_1, \lambda_2)}^1$ | |
| $(\lambda_1, \lambda_2) = (2i, 2j)$ | $\frac{4(1+(-1)^{i+j})}{(i+j+1)(2j+1)}$ |
| $(\lambda_1, \lambda_2) = (2i, 2j+1)$ | $\frac{-4(1+(-1)^j)}{(2i+2j+3)(j+1)}$ |
| Otherwise | 0 |
| $\mathcal{B}_{(\lambda_1, \lambda_2)}^s$ | |
| $(\lambda_1, \lambda_2) = (0, 0)$ | 8 |
| $(\lambda_1, \lambda_2) = (2i, 1)$ | $\frac{16}{(2i+3)(2i+1)}$ |
| $(\lambda_1, \lambda_2) = (2i, 2j), i+j \neq 0$ | $\frac{8(1+(2i+2j+1)(-1)^{i+j+1})}{ \text{Stab}_W(2i+1, 2j) (i+j+1)(i+j)(4j^2-1)}$ |
| $(\lambda_1, \lambda_2) = (2i, 2j+1), j \neq 0$ | $\frac{8(-1+(2j+1)(-1)^j)}{(2i+2j+3)(2i+2j+1)(j+1)}$ |
| Otherwise | 0 |
| $\mathcal{B}_{(\lambda_1, \lambda_2)}^l$ | |
| $(\lambda_1, \lambda_2) = (0, 0)$ | 8 |
| $(\lambda_1, \lambda_2) = (2i, 0), i \neq 0$ | $8 \left[\frac{2i+1+(-1)^{i+1}}{2i(i+1)} + \frac{1}{2i+1} \right]$ |
| $(\lambda_1, \lambda_2) = (2i, 2j), j \neq 0$ | $\frac{4}{ \text{Stab}_W(2i, 2j+1) } \left[\frac{2i+2j+1+(-1)^{i+j+1}}{(i+j+1)(2j+1)(i+j)} + \frac{2j+1+(-1)^{j+1}}{(2i+2j+1)(j+1)} \right]$ |
| $(\lambda_1, \lambda_2) = (2i, 2j+1)$ | $\frac{16}{ \text{Stab}_W(2i, 2j+2) } \left[\frac{(-1+(-1)^{j+1})(i+j+1)}{(2i+2j+3)(j+1)(2i+2j+1)} + \frac{(-1+(-1)^{i+j+1})(j+1)}{(i+j+1)(2j+3)(2j+1)} \right]$ |
| Otherwise | 0 |

TABLE 4. Values of \mathcal{B}_λ^t , given by (13), for $t \in \{0, 1, s, l\}$ and $\lambda = \lambda_1\omega_1 + \lambda_2\omega_2$. The indices i, j are non-negative integers.

$$g_5(y_1, y_2) = \begin{cases} r_5 & \text{if } y_1^2 + (y_2 + 1.5)^2 \leq 1, \\ 0 & \text{otherwise.} \end{cases}$$

graphs of of the absolute value of the difference $|1 - \mathcal{I}_M^t(g_i)|$.

Each value of $r_i \in \mathbb{R}$ is set to satisfy the normalization condition $\int_T g_i(y) dy = 1$. We compute the approximations

$$\mathcal{I}_M^t(g_i) = \sum_{\substack{\lambda \in P^+ \\ \langle \lambda, \eta \rangle \leq M}} b_\lambda^t \mathcal{B}_\lambda^t \tag{14}$$

of $\int_T g_i(y) dy = 1$ with the formula for \mathcal{B}_λ^t given by (13). Figs. 8 and 9 show for $t \in \{0, 1, s, l\}$ the

5. CONCLUDING REMARKS

- (1.) Establishing the explicit connection between the Jacobi and Macdonald polynomials and the Weyl group orbit functions in Section 2.3 forms a crucial step for generalizing known cubature formulas to the entire class of the Jacobi polynomials.
- (2.) Numerical tests results in Figs. 5 and 8 indicate in general excellent convergence rates of the developed cubature rules including their Clenshaw-Curtis

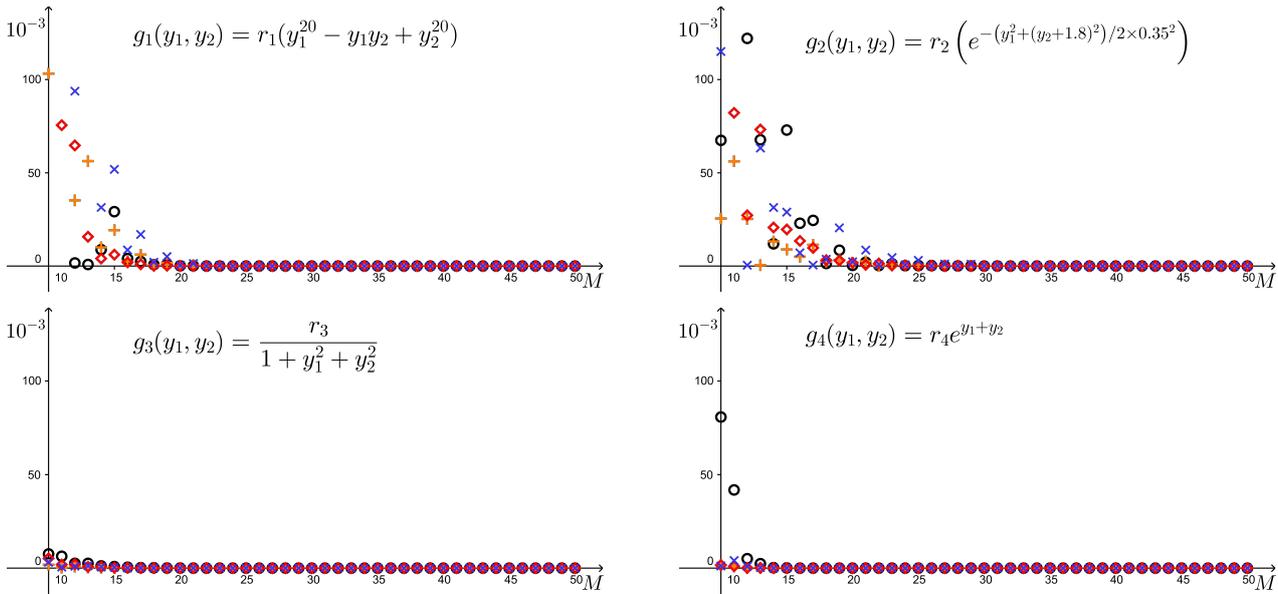


FIGURE 8. The graphs of error values $|1 - \mathcal{I}_M^t(g_i)|$ of the integral $\int_T g_i(y) dy = 1, i = 1, \dots, 4$ and its approximations $\mathcal{I}_M^t(g_i), M = 10, 11, \dots, 50$ given by (14). The values for $t = 0, 1, s, l$ are depicted as circles, “+”, diamonds and “x”, respectively.

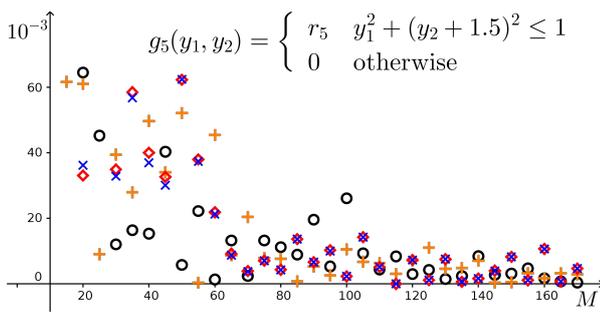


FIGURE 9. The graphs of error values $|1 - \mathcal{I}_M^t(g_5)|$ of the integral $\int_T g_5(y) dy = 1$ and its approximations $\mathcal{I}_M^t(g_5), M = 10, 15, 20, \dots, 170$ given by (14). The values for $t = 0, 1, s, l$ are depicted as circles, “+”, diamonds and “x”, respectively.

versions for the case C_2 . The convergence rate of the multidimensional step-functions, even though less uniform, still appears to be very good. Developing similar methods for the two-variable case G_2 and extending the rules to higher dimensions poses an open problem.

- (3.) The hyperinterpolation methods [3, 25, 34, 35] are among the tools which directly use cubature rules. For the standard cubature rules of the Weyl group orbit functions, several tests with very good results are also performed in [15]. Developing and testing hyperinterpolation methods for the presented cubature rules merits further study.
- (4.) The present work demonstrates wide variety of possibilities of constructing the cubature rules in the orbit functions setting. Comparison of the developed methods is necessary for establishing range of their viable applications. Especially, comparison

of the Gauss and Clenshaw-Curtis cubature methods, similar to [37], regarding their efficiency, speed, model function and integration domain dependence merits further research.

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